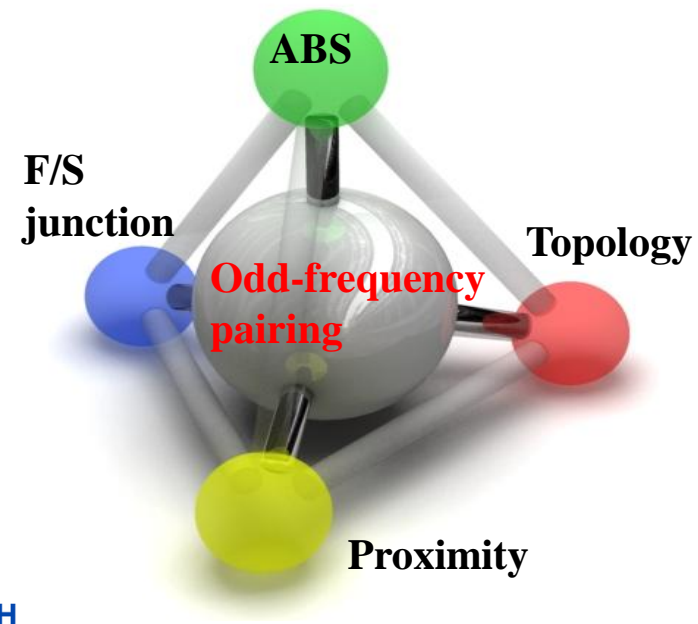


# Structures Superconductor – Topological Insulator: Theory and Experiment

Alexander Golubov

*University of Twente, The Netherlands*

*Moscow Institute of Physics and Technology, Russia*



**MESA+**  
INSTITUTE FOR NANOTECHNOLOGY

**MIPT**  
MOSCOW INSTITUTE  
OF PHYSICS AND TECHNOLOGY

  
SPINTECH

# *Collaborators*

**V.S. Stolyarov, *Moscow Institute of Physics and Technology***

**D. Rodichev, *Universite Pierre et Marie Curie, Paris***

**V.M. Vinokur, *Argonne National Laboratory***

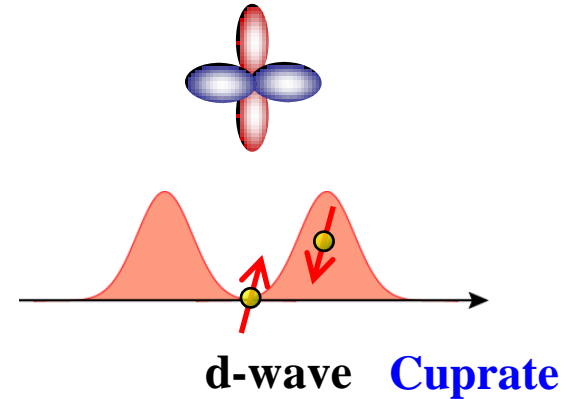
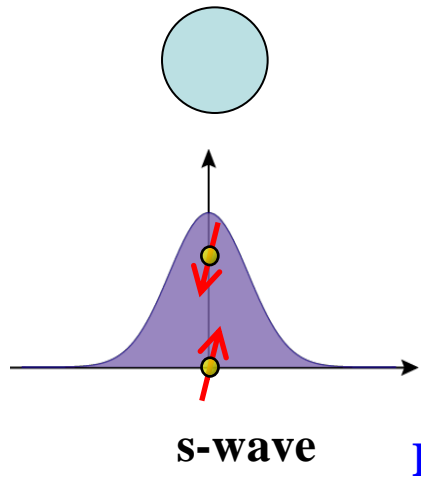
**M.Yu. Kupriyanov, *Moscow State University***

# Content

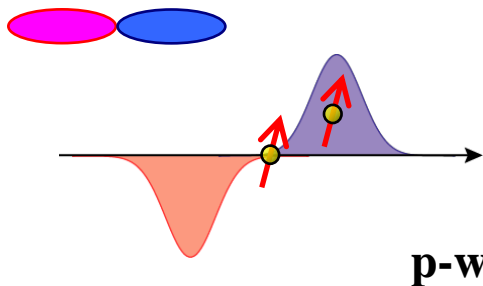
- (1) Unconventional superconductivity
- (2) Andreev bound states in unconventional junctions
- (3) Search of p-wave superconductivity in topological materials

# Classification of Symmetry of Cooper pair

**Spin-singlet Cooper pair**  $\longrightarrow$  **Even Parity**



**Spin-triplet Cooper pair**  $\longrightarrow$  **Odd Parity**



$^3\text{He}$   $\text{Sr}_2\text{RuO}_4$

# Topological superconductors

Corresponding edge state in the world of superconductivity (superfluid)

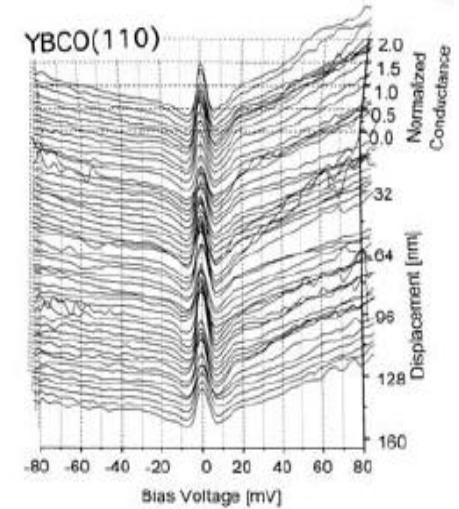
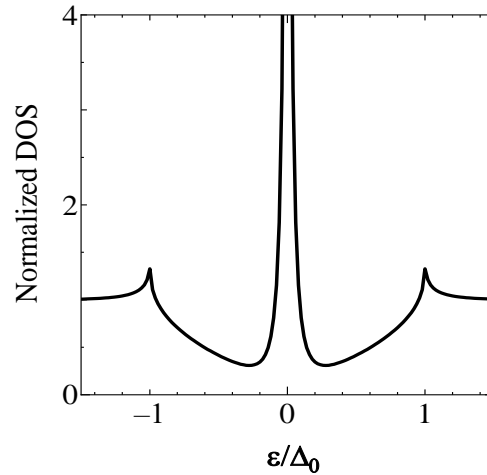
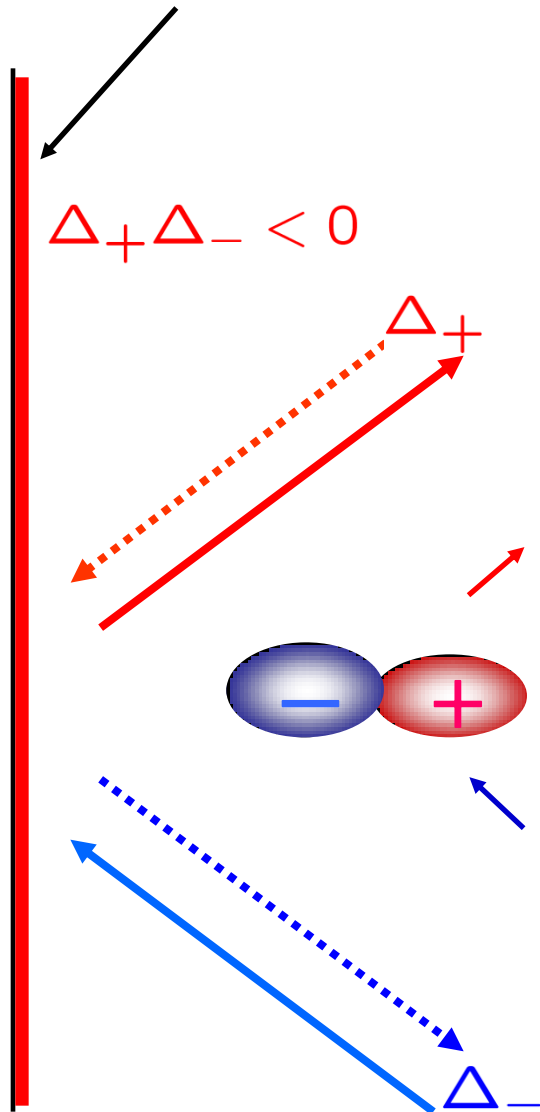


**Andreev bound state** generated at the surface of superconductors by the bulk-edge correspondence

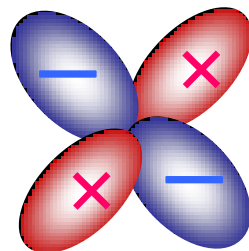
X. L. Qi, S.C. Zhang, PRL 102, 187001 (2009), Schnyder et. al, PRB 78 195125 (2008)

M. Sato and Fujimoto, PRB 79 094504 (2009), Y. Tanaka et. al, PRB 79 060505 (2009)

# Classical example of Andreev bound state in unconventional superconductors (high Tc cuprate)



Local density of state has a zero energy peak. (Sign change of the pair potential at the interface)




Tanaka Kashiwaya PRL 74 3451 (1995); Hu(1994)  
Kashiwaya, Tanaka, Rep. Prog. Phys. 63 1641 (2000)

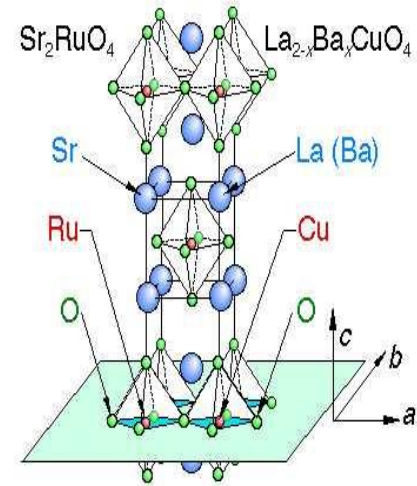
Buchholz, Zwicky(1981) Hara Nagai(1986)  
Hu(1994), M. Sato(2011)

# Chiral superconductor $\text{Sr}_2\text{RuO}_4$

Edge surface current



  $p_x + ip_y$

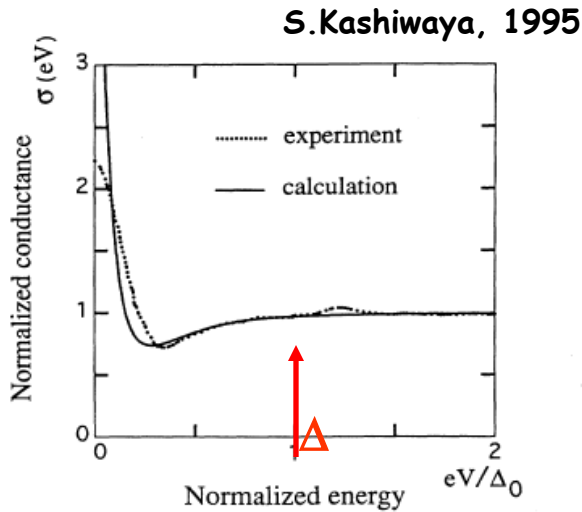


Similar structure to cuprate



Maeno (1994)

# Tunneling spectrum in cuprate and Ruthenate



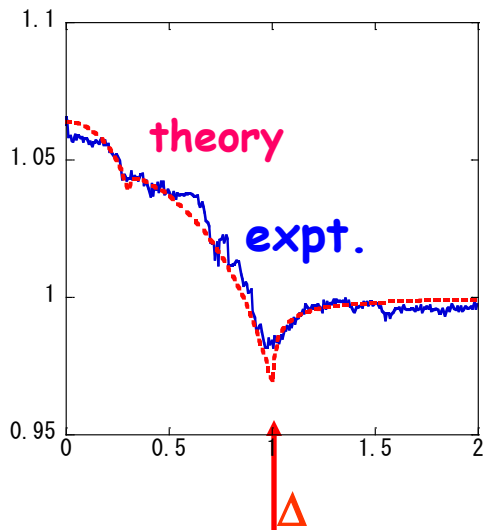
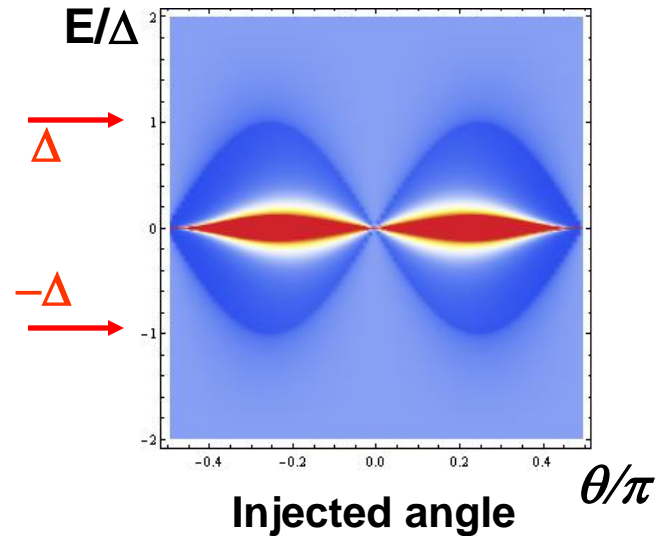
**YBCO(110)**

$d_{x^2-y^2}$ -wave

**Zero energy state**

**Dispersionless bound state**

**Angle resolved conductance**

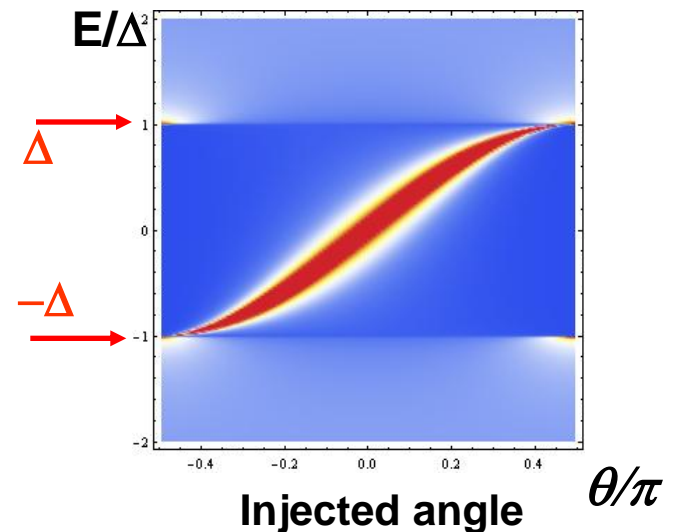


**Sr<sub>2</sub>RuO<sub>4</sub>**

chiral p-wave

**Chiral edge state**

**Broad peak due to linear dispersion**



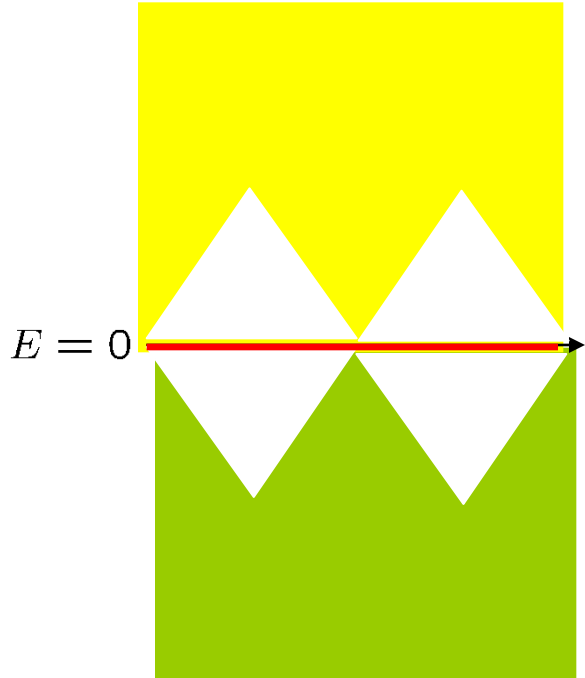
Phys. Rev. Lett. 107, 077003 (2011)



# Feature of the Andreev bound states

Non-centrosymmetric superconductor (NCS)

$d_{xy}$ -wave

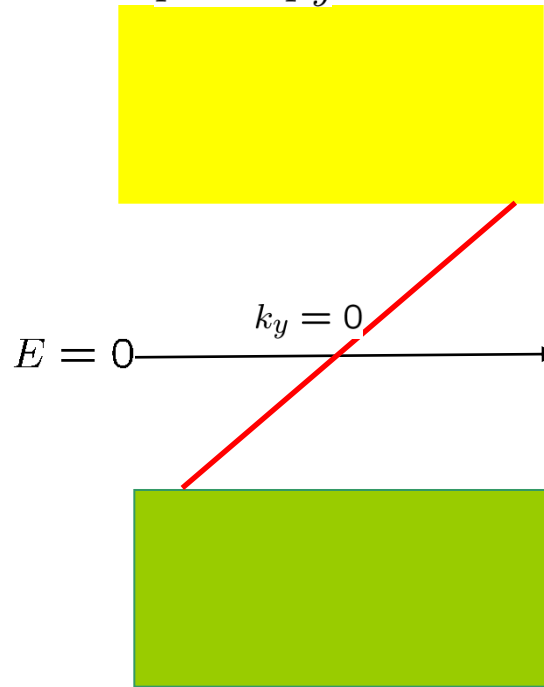


Hu(94)

Tanaka Kashiwaya (95)

**Flat**

Chiral p-wave  
 $p_x + ip_y$  -wave

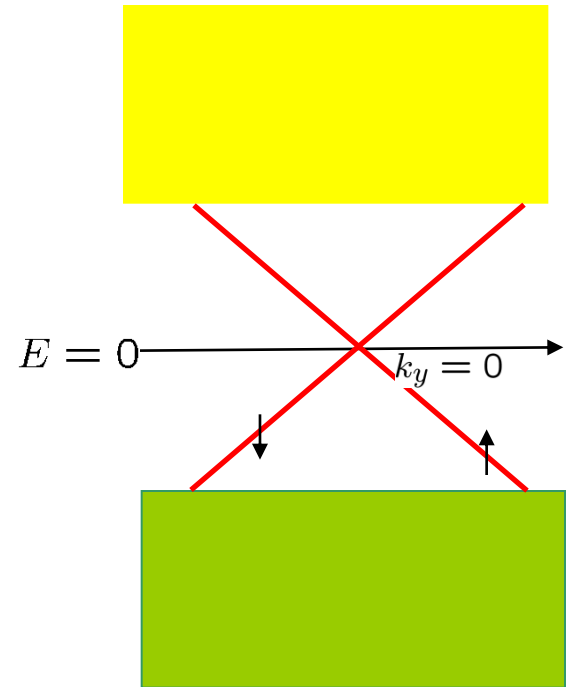


Tanaka Kashiwaya (97)

Sigrist Honerkamp (98)

**Chiral**

NCS (Helical)  
 $p+s$  -wave



Iniotakis (07)

Eschrig(08)

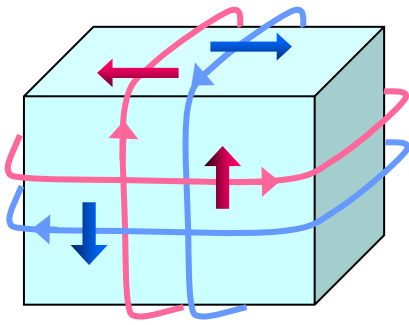
Tanaka (09)

**Helical**

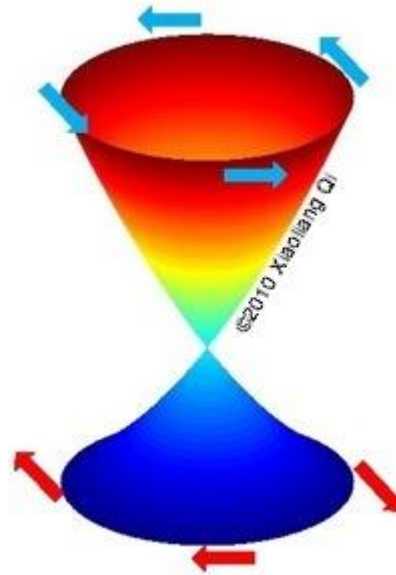
# **Superconducting junctions with topological materials**

# Topological insulator 3D systems

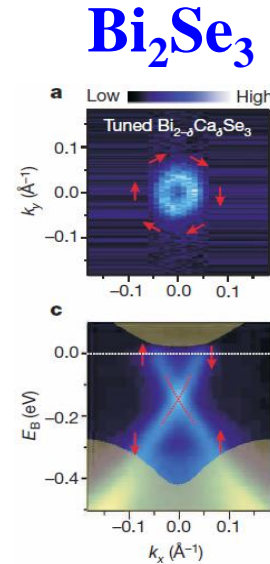
Spontaneous  
Spin current



Fu Kane 2007  
J. Moore 2007

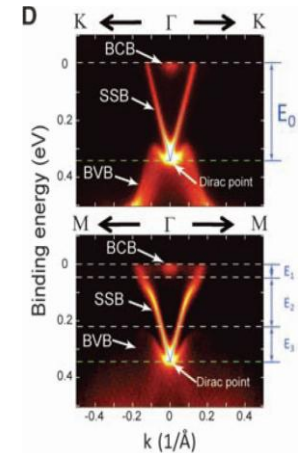


Surface  
Dirac Cone



Hsieh et al., Nature (2009)

**Bi<sub>2</sub>Te<sub>3</sub>**

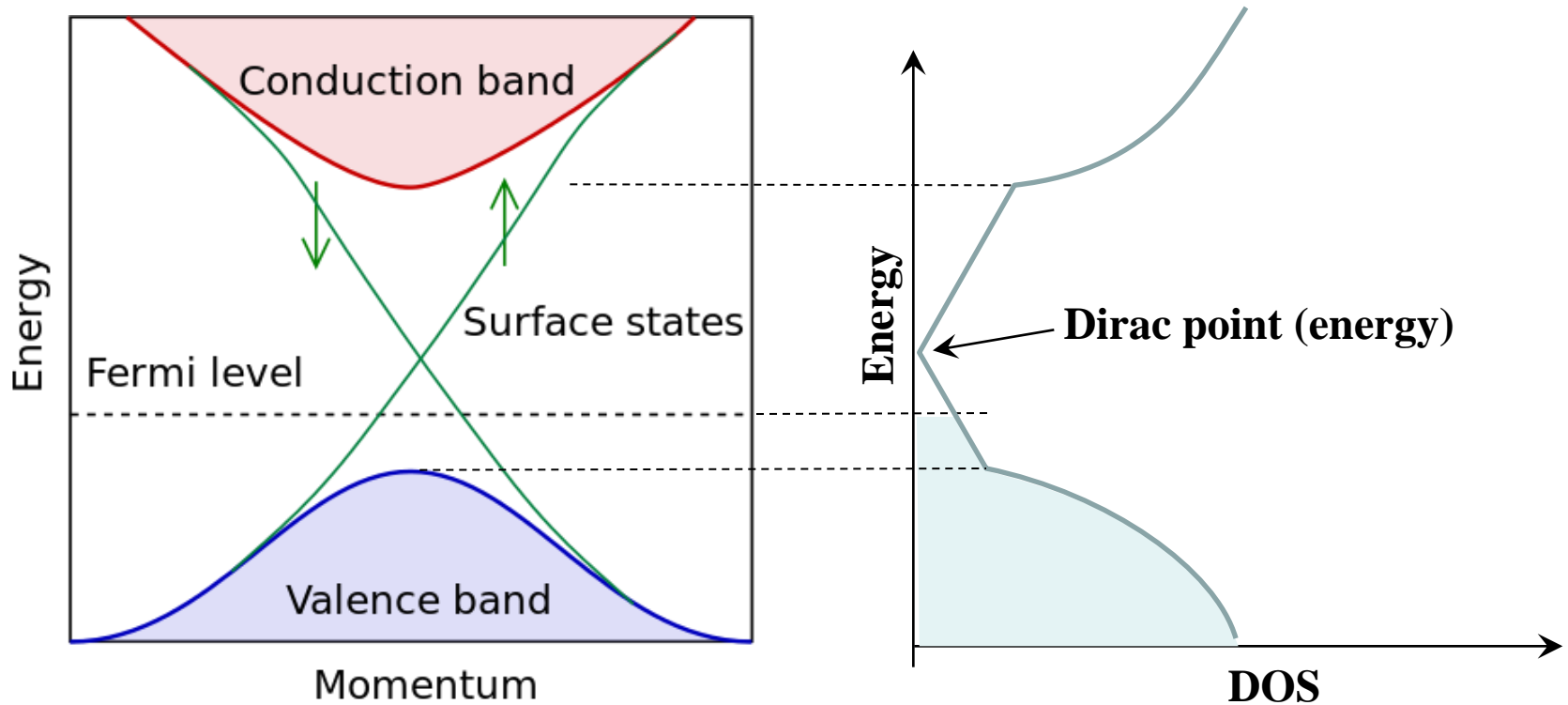


Quantum number  $Z_2$  which specifies topological insulator

(Product of parities of wave function  
at time reversal invariant points in  $k$ -space)

Fu Kane 2007

## 3D Topological insulators: Ideal case



[https://en.wikipedia.org/wiki/Topological\\_insulator](https://en.wikipedia.org/wiki/Topological_insulator)

**Motivation:** high metallicity, forbidden backscattering, specific momentum-spin relations, possible induction of topological (p-wave) superconductivity, etc.

# 2D case: *p*-wave pairing symmetry

Pair wave function:

orbital x **spin** x **frequency** = asymmetric

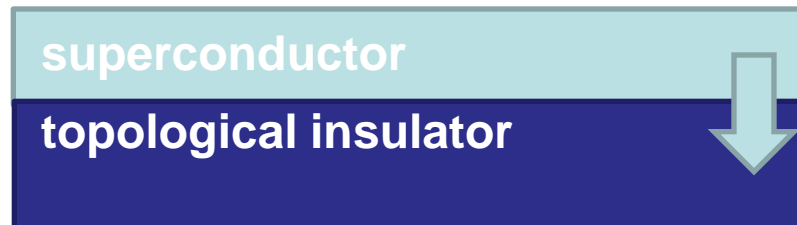
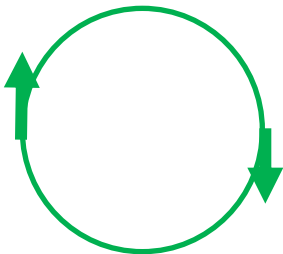
Simplest superconductor (e.g. Nb or Al):

$|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle$  s-wave, **singlet**, **even freq.**

Superconductivity in topological insulator:

$$|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \xrightarrow{\text{X}} \frac{1}{2} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) + \frac{1}{2} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$

50% s-wave, **triplet**, **even freq.** + 50% p-wave, **triplet**, **even freq.**



proximity  
effect

# 2D-case: *p*-wave pairing symmetry

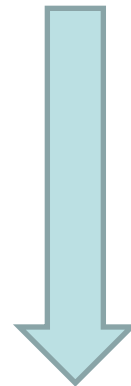
Superconductivity in topological insulator:

50% s-wave, **singlet**, **even freq.** + 50% p-wave, **triplet**, **even freq.**

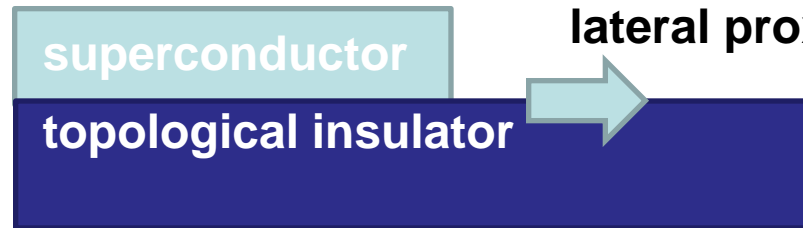


breaking *time-reversal* symmetry

→ Towards 100% p-wave, **triplet**, **even freq.**



breaking *translational* symmetry

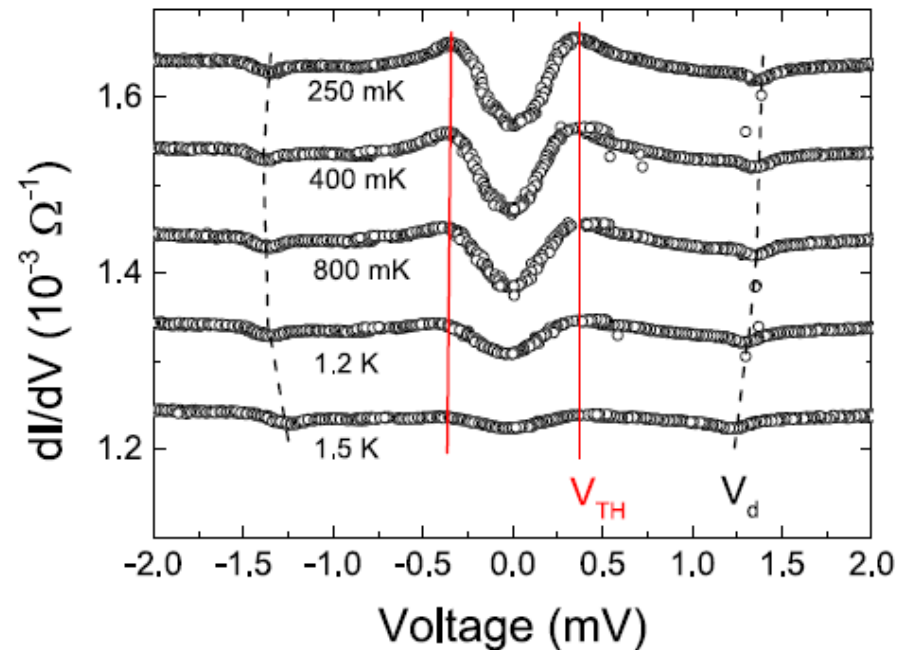
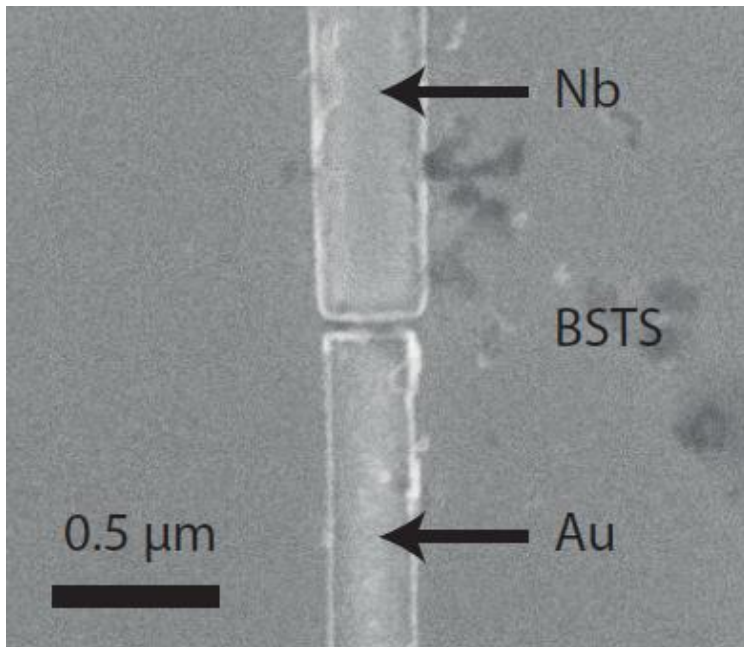
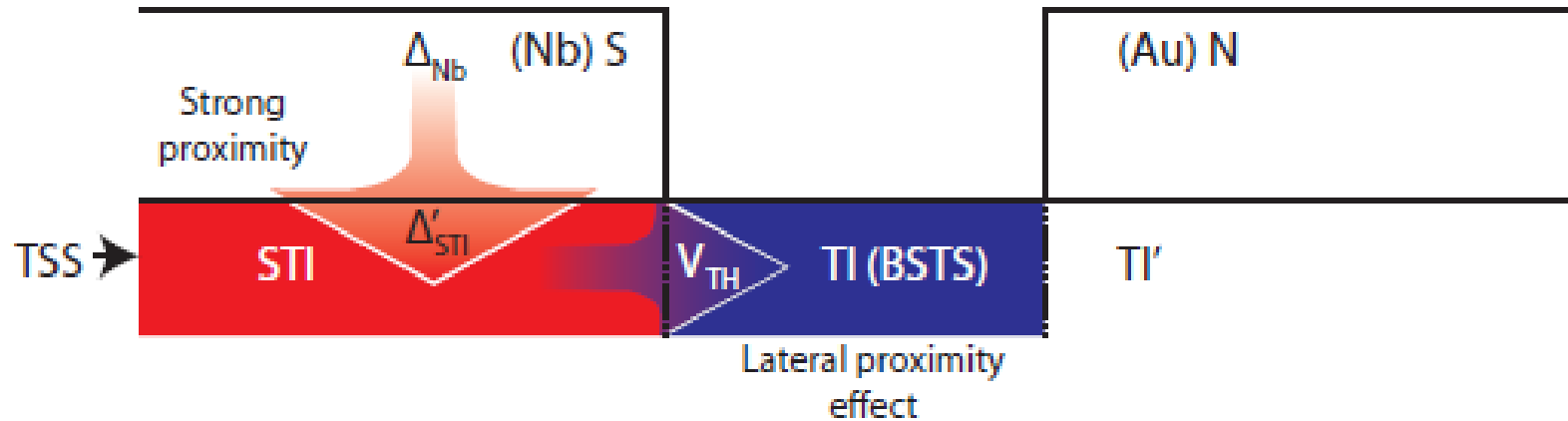


→ 100% s-wave, **triplet**, **odd freq.**

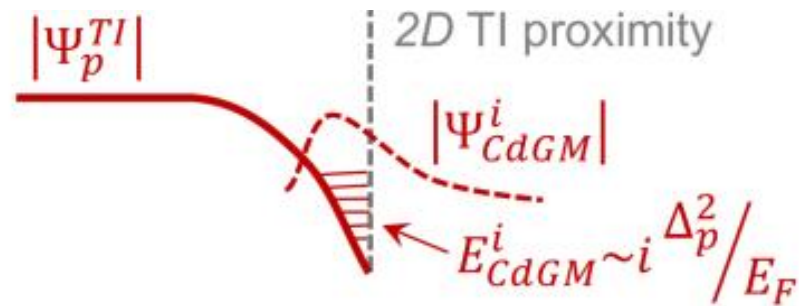
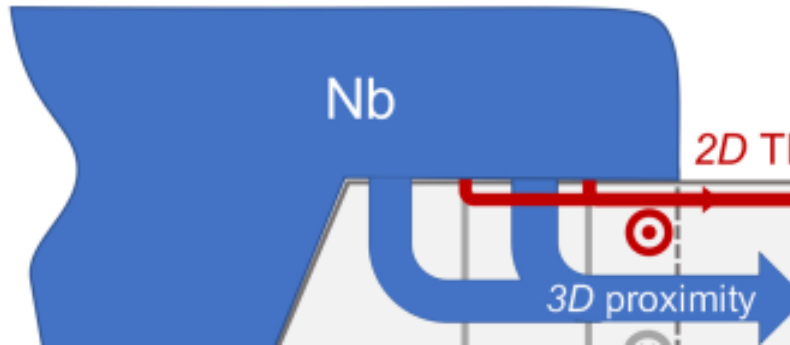
# dl/dV proximity spectroscopy, experiment (University of Twente).

Diffusive regime: *no evidence of p-wave state*

(a)



# What can we expect in the ballistic regime ?



*Normal reflections occur near the edge of the S – layer =>  
suppression of the p-wave order parameter near the edge =>  
fine structure of the energy levels*

*(similar to the Caroli - de Gennes – Matricon levels in a vortex core)*



Michael Stone PRB 54, 13 222 (1996)

## *Geometric optics limit of Andreev reflection*

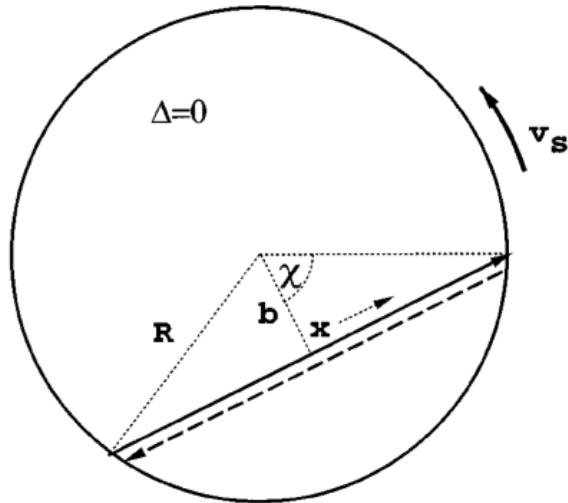


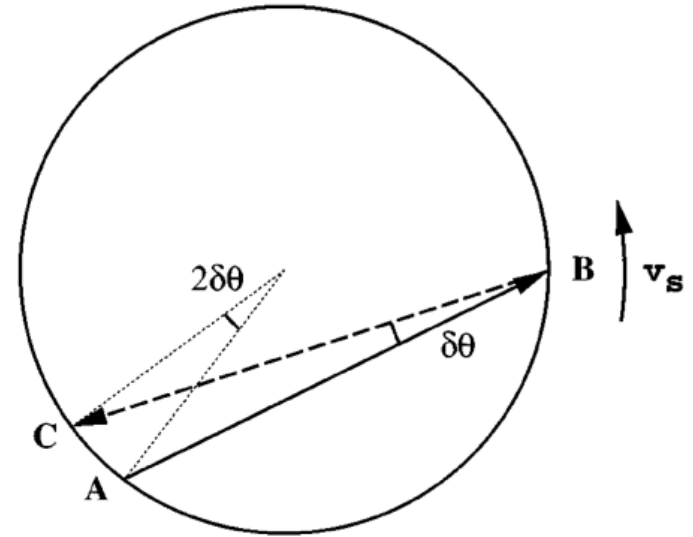
FIG. 1. A bound state with an electron (solid arrow) being Andreev reflected as a hole (dashed arrow).

### *Caroli - de Gennes – Matricon levels in a vortex core*

$$\epsilon_n(l) = \frac{v_f}{4R \sin \chi} [2\chi + 2\pi(n + \frac{1}{2})]$$

$$\epsilon_{-1}(l) = -\omega_0 l, \quad \omega_0 = \frac{1}{2mR^2}.$$

$$R \approx v_f / \Delta \quad \Rightarrow \quad \epsilon \sim \Delta^2 / \epsilon_F$$



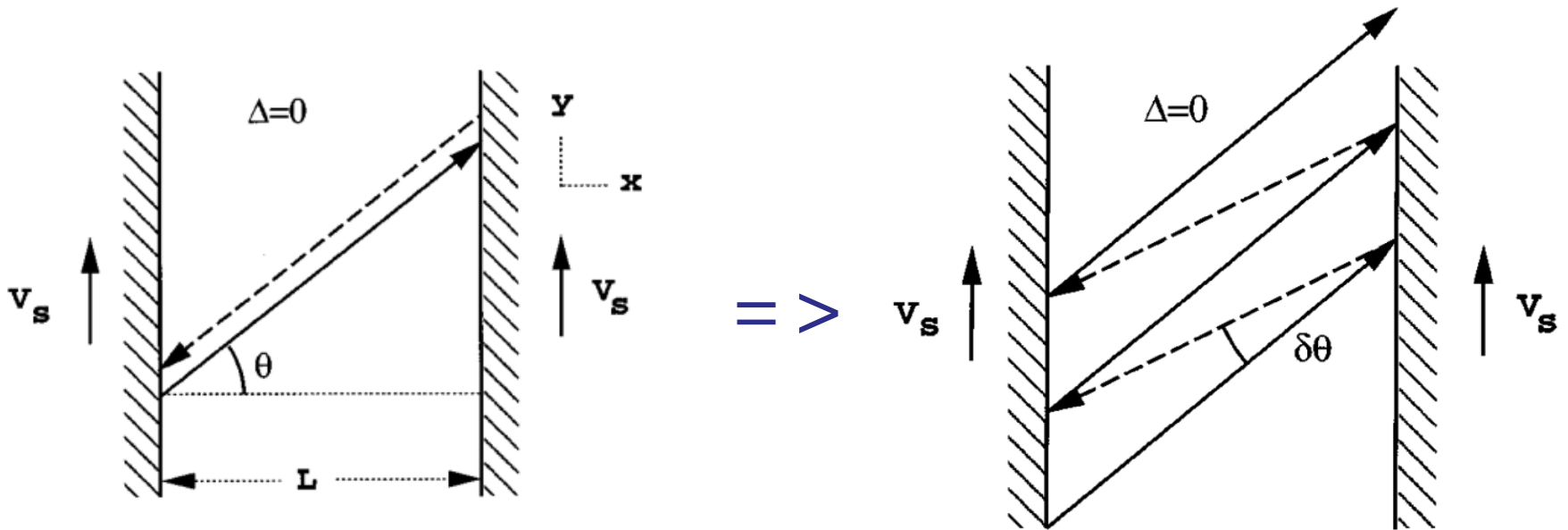
**Physical mechanism:**  
*loss of retroreflection property  
of the Andreev reflection process  
(beyond the quasiclassical approximation)*

$$\delta\theta \sim v_s / v_f$$

# SN interfaces in the presence of superflow

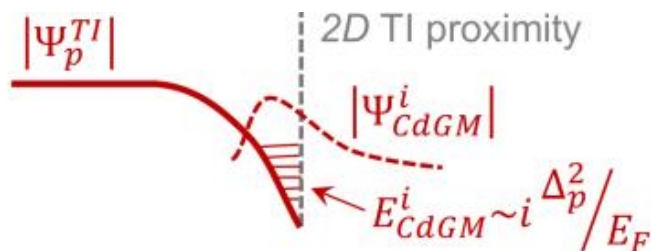
M PRB .54, 13 222 (1996)

*loss of retroreflection property of the Andreev process  
(beyond the quasiclassical approximation)*



$$\delta\theta \sim v_s/v_f$$

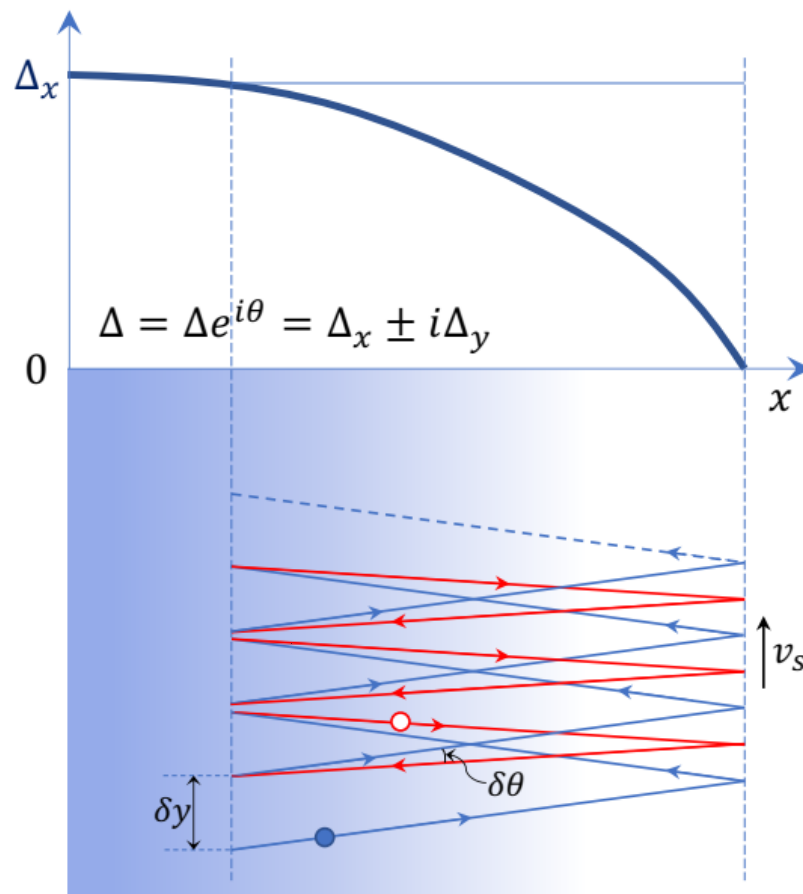
# Application of the *geometric optics regime of Andreev reflection* to the topological junction



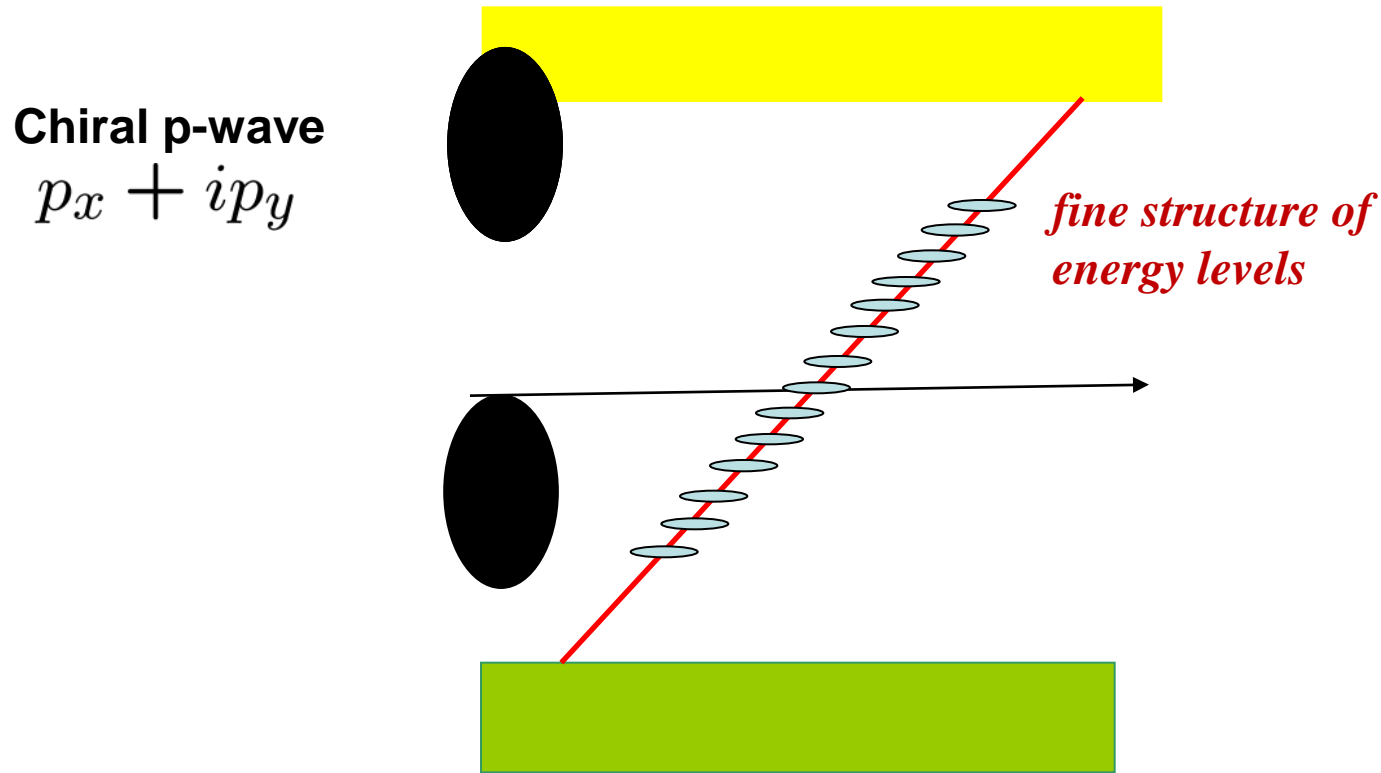
$$\epsilon_n(l) = \frac{v_f}{4R \sin\chi} [2\chi + 2\pi(n + \frac{1}{2})]$$

In the p-wave case:  $2\chi = \pi + 2\delta\theta$ ,  $R \sim \xi$

$$\delta\theta \sim v_s/v_f \Rightarrow \epsilon \sim \Delta^2/\epsilon_F$$



*The result: splitting of the linear spectral branch into discrete energy levels*



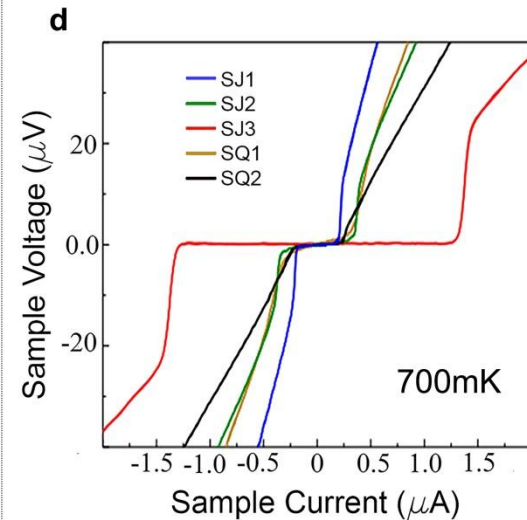
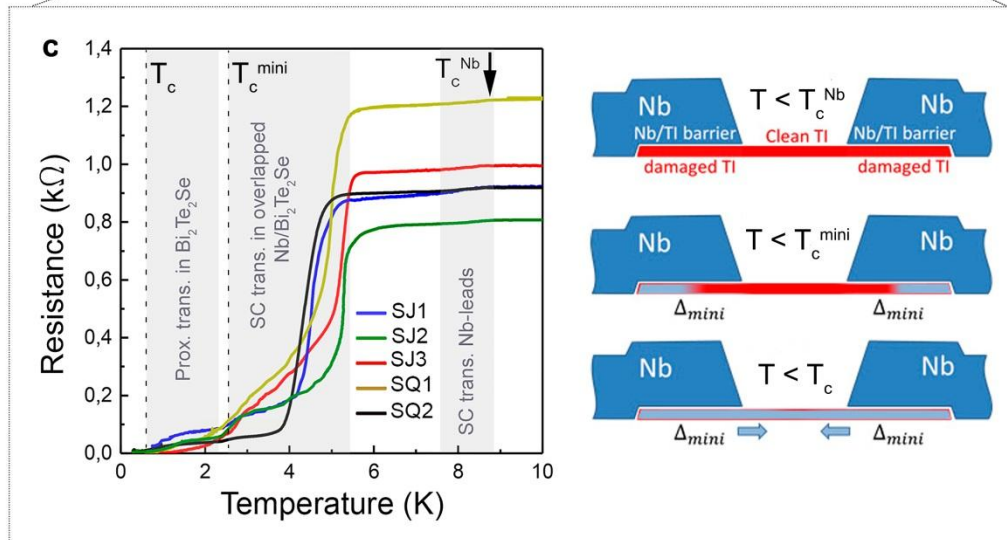
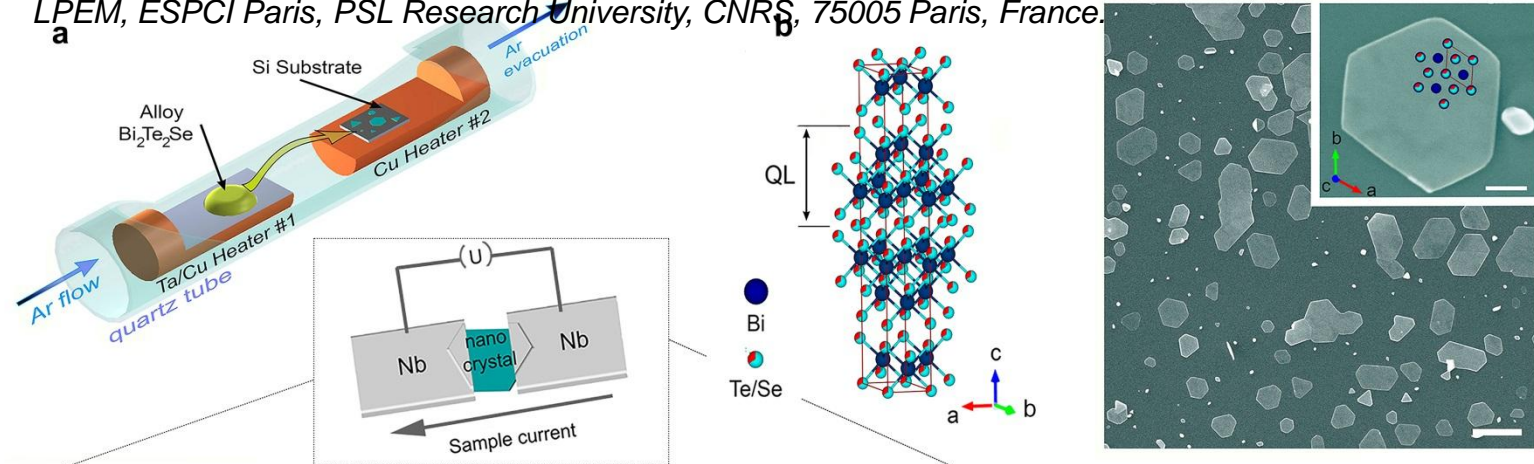
*Fine structure should lead to small-scale periodicity of junction properties, e.g. to short-periodic oscillations in a magnetic field*

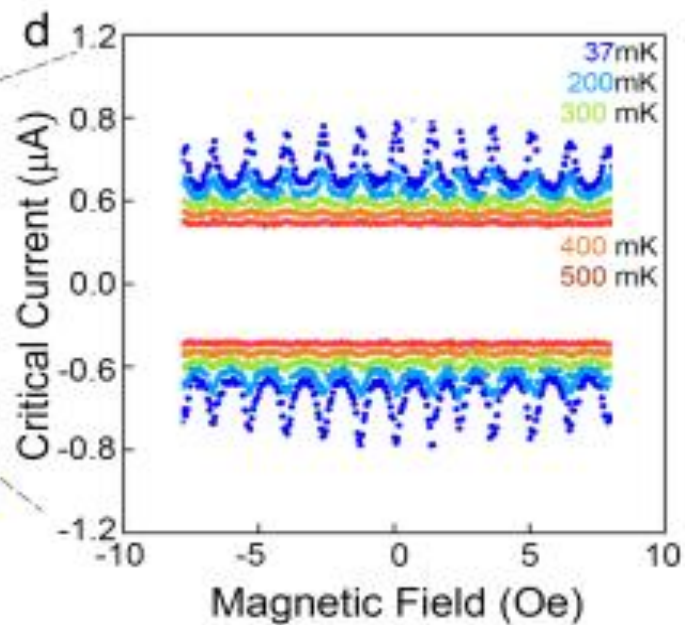
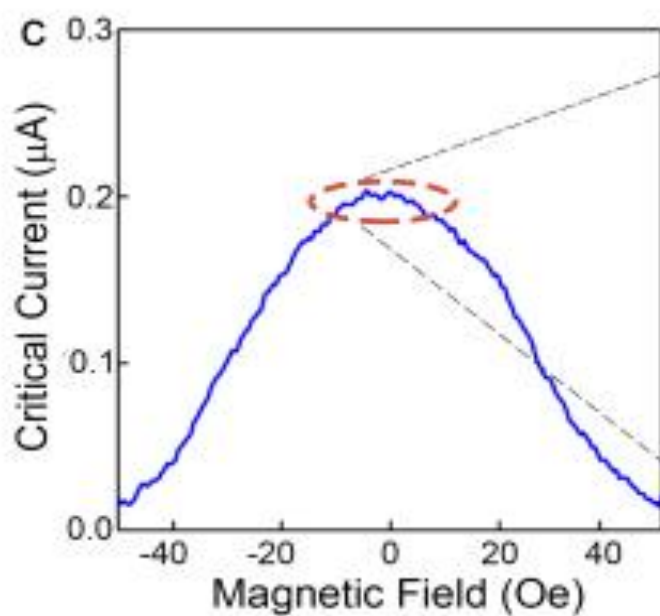
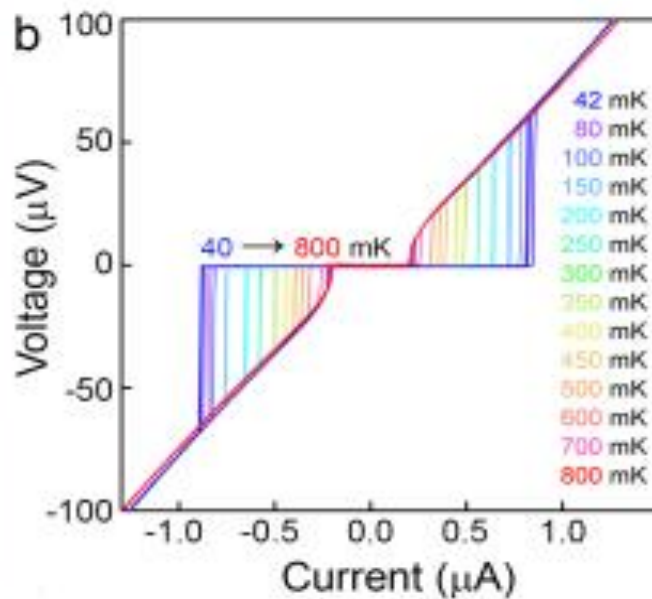
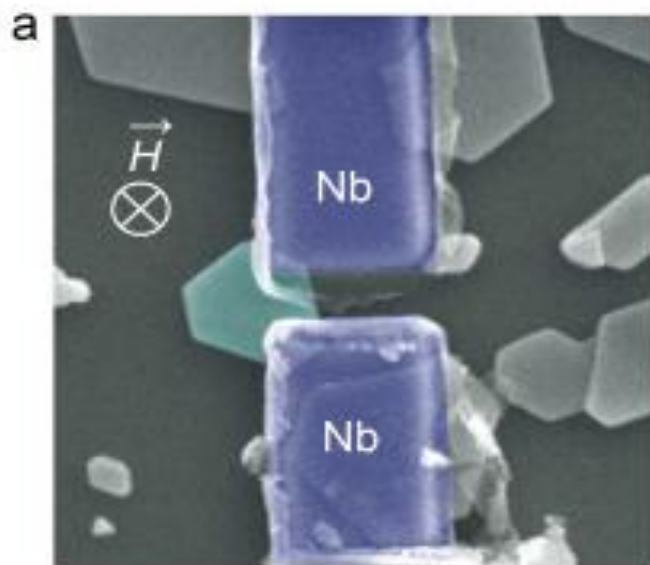
# Recent experiment: Nb - Bi<sub>2</sub>Te<sub>2.3</sub>Se<sub>0.7</sub> – Nb Josephson junctions

V.S. Stolyarov, *Communication Materials* (2020)

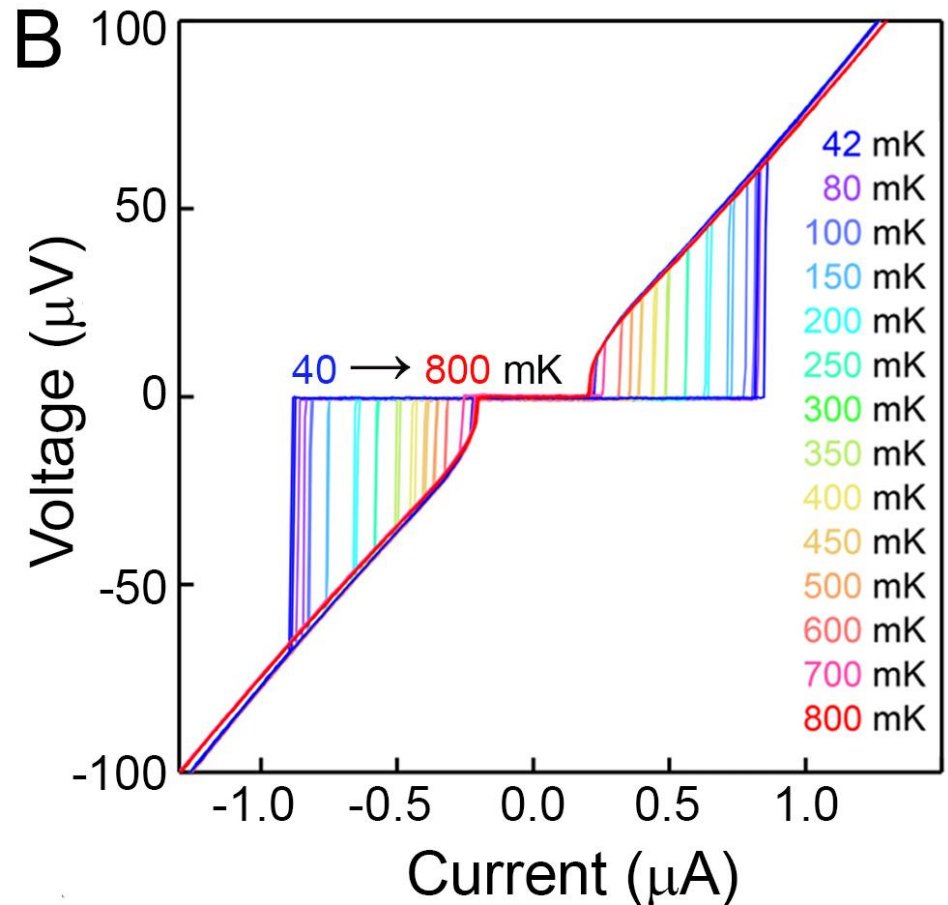
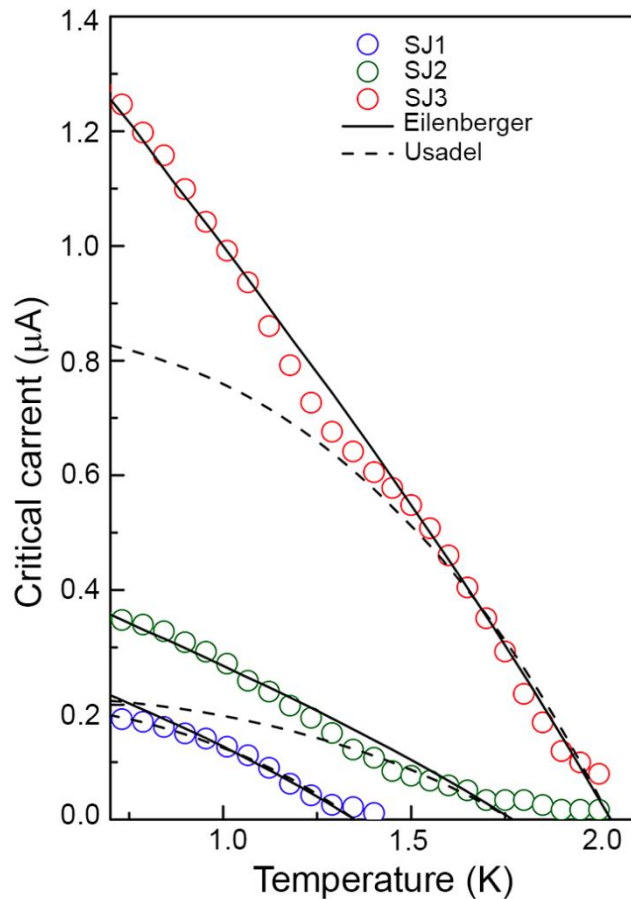
Moscow Institute of Physics and Technology, Russia

LPEM, ESPCI Paris, PSL Research University, CNRS, 75005 Paris, France.





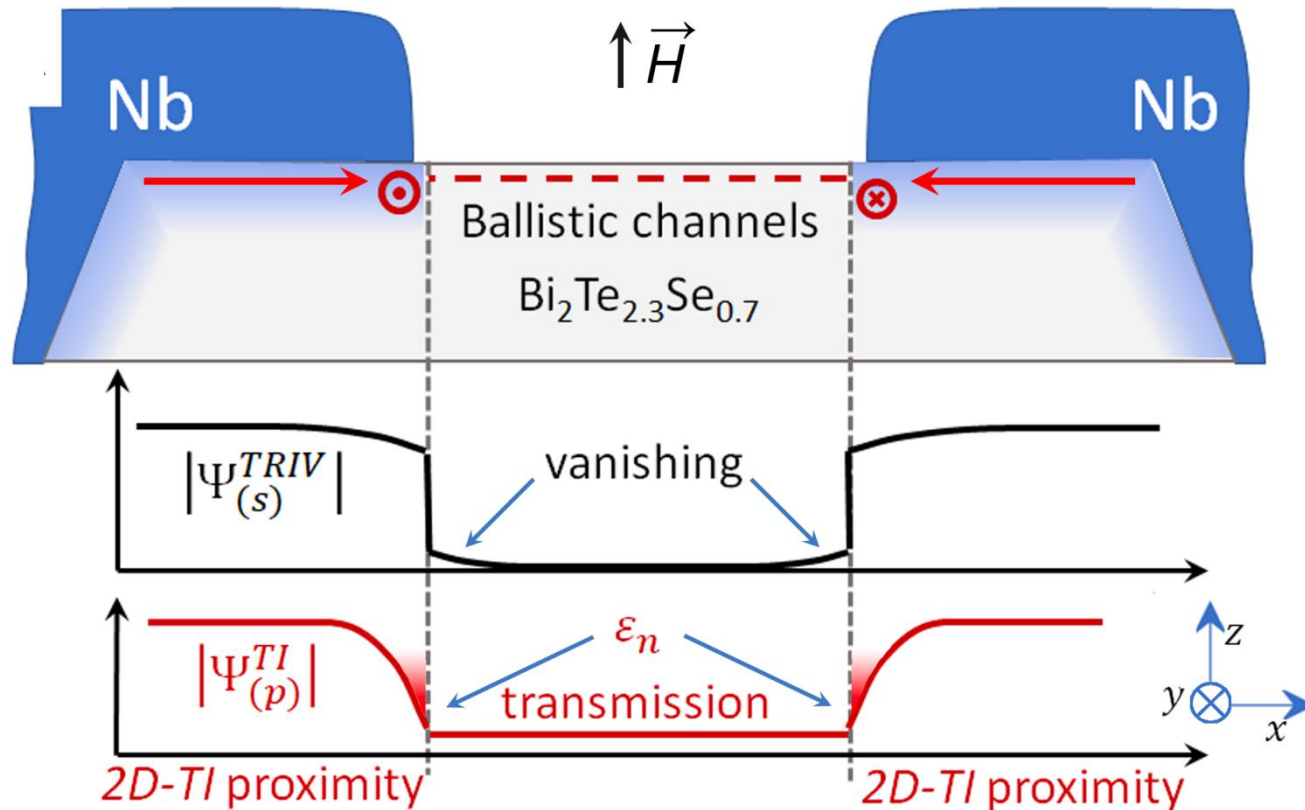
# Josephson phenomena in Nb-Bi<sub>2</sub>Te<sub>2.3</sub>Se<sub>0.7</sub>-Nb S-TI-S junctions



*V. S. Stolyarov, et al. Communications Materials 1, 38 (2020)*

**Transmittance by  $\sim 10$  ballistic surface channels; no bulk**

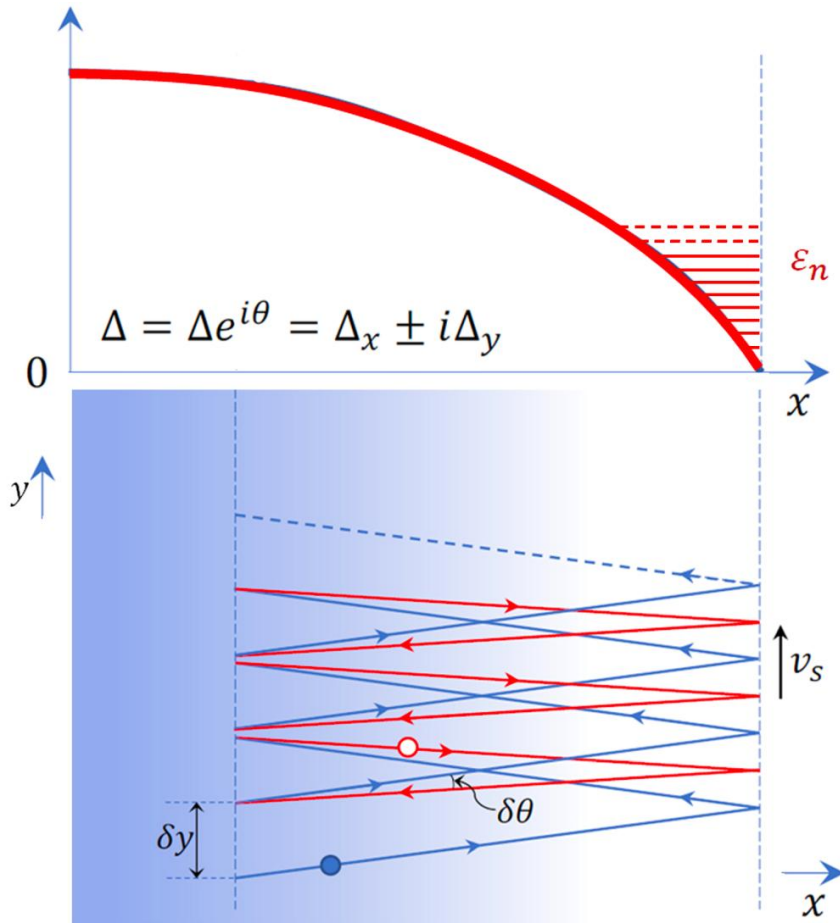
# Josephson phenomena in Nb-Bi<sub>2</sub>Te<sub>2.3</sub>Se<sub>0.7</sub>-Nb S-TI-S junctions



**Hypothesis: two-component superconducting correlations; ballistic topological channels carry only  $p$  – component**



# Josephson phenomena in Nb-Bi<sub>2</sub>Te<sub>2.3</sub>Se<sub>0.7</sub>-Nb S-TI-S junctions



Andreev bound states with dispersion:

$$\varepsilon = \Delta \sin(k_y y)$$

$$\delta\theta \sim v_s/v_F \sim \Delta/\varepsilon_F$$

mid-gap state at  $k_y = 0$

$$\varepsilon_n \sim \Delta(v_s/v_F)n \sim (\Delta^2/\varepsilon_F)n$$

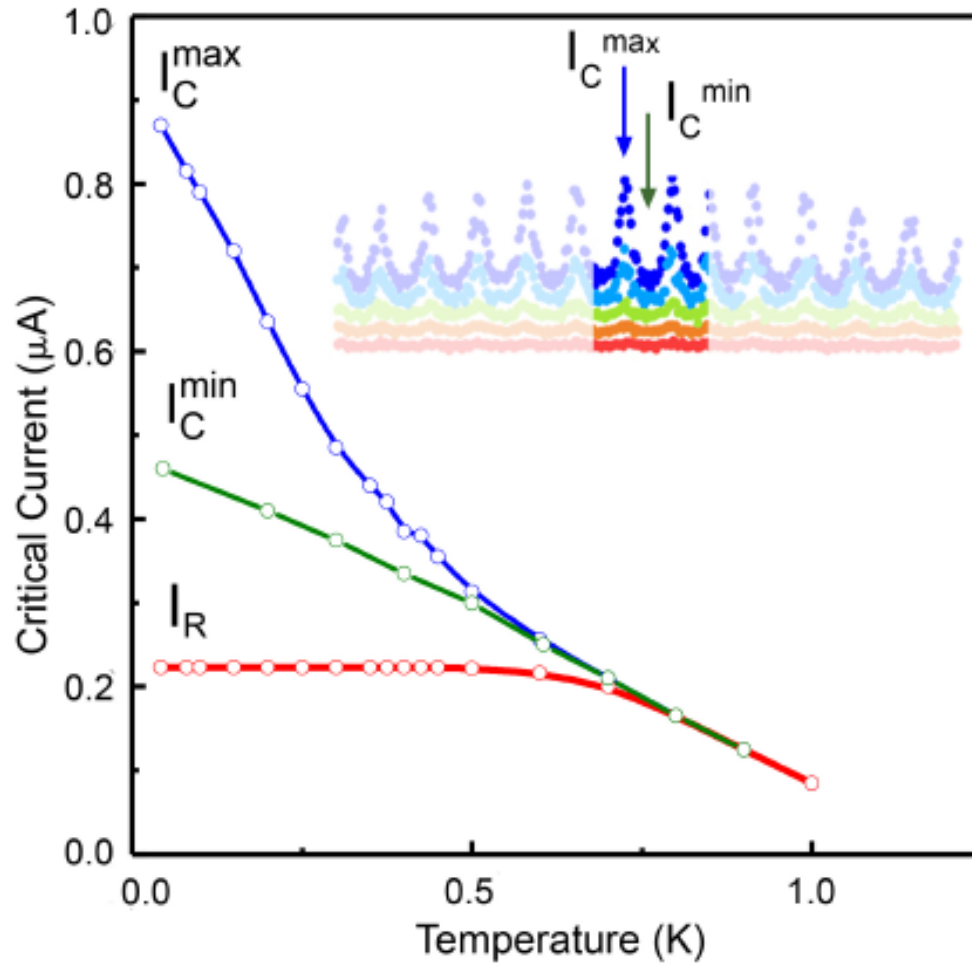
i.e. similar to Caroli-Matricon-deGennes states in the vortex core!

In the case of Bi<sub>2</sub>Te<sub>2.3</sub>Se<sub>0.7</sub> ( $\Delta \approx 0.3\text{meV}$ ,  $\varepsilon_F \approx 0.1\text{eV}$ ),

$$\frac{\Delta^2}{\varepsilon_F} \sim 1 \mu\text{eV}$$

Sort of Fabry-Perot interferometer with phase-shifting mirrors.  
Magnetic field creates an extra phase shift.

# Josephson phenomena in Nb-Bi<sub>2</sub>Te<sub>2.3</sub>Se<sub>0.7</sub>-Nb S-TI-S junctions

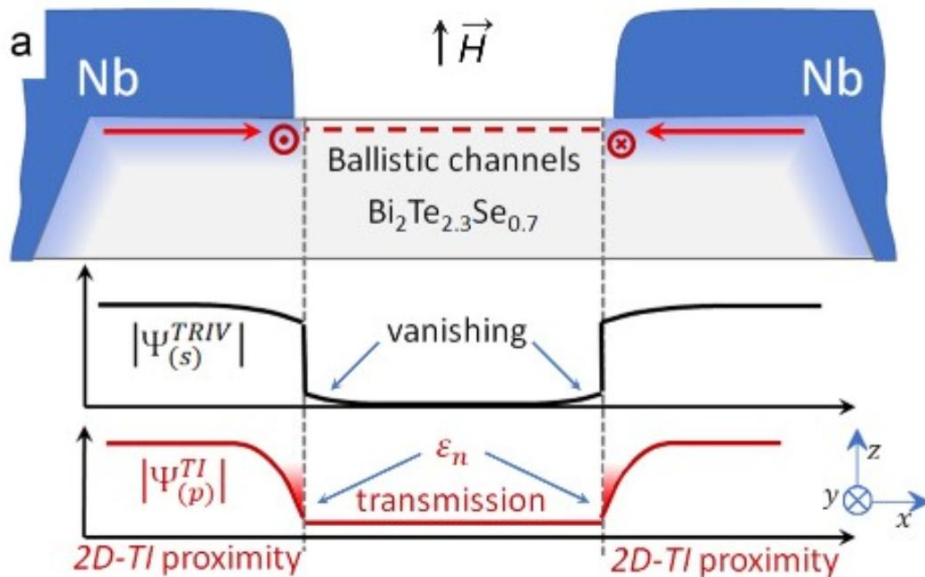


**Observed strong temperature dependence is expected :  $\sim \Delta^2/T$**

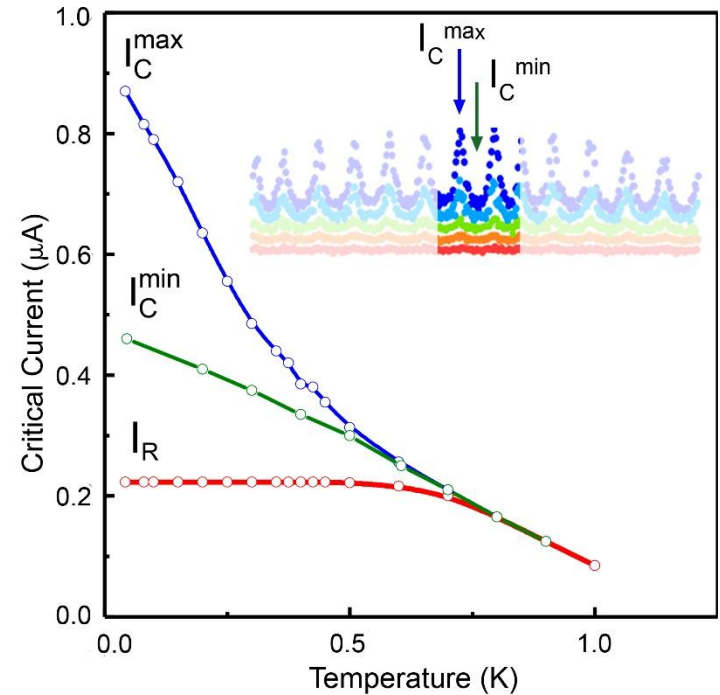
*Tanaka, Y. & Kashiwaya, S. Phys. Rev. B, 56, 892912 (1997)*

*Ilichev, E. et al. Phys. Rev. Lett. 86, 5369 (2001).*

*The expected spatial variation of s-wave and p-wave correlations inside the nanocrystal*



*The evolution of the maximum and minimum critical currents. Inset:  $T = 37$  mK*



*The observed oscillations reflect the resonant contribution from the Andreev levels to  $I_C$*

*The data provide the evidence of the p-wave topological nature of the proximity-induced superconductivity at the TI surface*

# *Summary*

- (1) Midgap Andreev bound states – key property of **unconventional superconductivity**
- (2) Signatures of p-wave superconductivity in diffusive systems:  
**anomalous proximity effect in DN/spin-triplet p-wave junction**  
due to generation of the odd-frequency pairing state
- (3) Signatures of p-wave superconductivity in ballistic systems:  
observed **fine structure of low-energy Andreev bound states**  
at the surface of *Bi<sub>2</sub>Te<sub>2.3</sub>Se<sub>0.7</sub>*